

# 1 Problem Set #1

1. Suppose that  $\rho = e^{-\beta H}$  is a thermal state in an ordinary Hilbert space, where  $H$  is any operator for which the resulting state is normalizable. Show that this state satisfies the KMS condition for any operators  $A, B$ :

$$\langle A(\tau)B \rangle_\rho = \langle BA(\tau + i\beta) \rangle_\rho, \quad (1)$$

where  $A(\tau)$  represents the time translation of the operator  $A$ , and  $\langle BA(z) \rangle_\rho$  is analytic in the strip  $0 \geq \text{Im}(z) \geq \beta$ .

(This exercise justifies the use of the KMS condition as the *definition* of a thermal state in more exotic contexts, such as type III von Neumann algebras.)

2. Consider the following properties of the von Neumann entropy described during the lecture:

- i) Positivity:  $S(\rho) \geq 0$ ,
- ii) Invariance under Unitaries:  $S(U\rho U^\dagger) = S(\rho)$ ,
- iii) Additivity under Tensor Product:  $S(\rho_A \otimes \rho_B) = S(\rho_A) + S(\rho_B)$ ,
- iv) Triangle Identity:  $S(A) + S(B) \geq S(AB) \geq |S(A) - S(B)|$ ,
- v) Continuous (for finite dimensional Hilbert spaces),
- vi) {Strong Subadditivity:  $S(AB) + S(BC) \geq S(ABC) + S(B)$ },
- vii) Concavity:  $S(\lambda\rho + (1 - \lambda)\sigma) \geq \lambda S(\rho) + (1 - \lambda)S(\sigma)$ ,
- viii) Chain Rule: If  $\rho = \otimes_i \lambda_i \rho_i$  (block diagonal), then  $S(\rho) = \langle S(\rho_i) \rangle - \sum_i \lambda_i \ln \lambda_i$ .

(a) Show that Strong Subadditivity implies Weak Monotonicity:  $S(AB) + S(BC) \geq S(A) + S(C)$  by assuming the existence of a 4th system  $D$  such that  $ABCD$  is pure, and using the fact that  $S(R) = S(\bar{R})$ .

(b) See how many of the properties you can prove on your own. (But don't try too hard for Strong Subadditivity, whose proof is very difficult!) For the proof of Araki-Lieb, try using the purifier trick described in (a). Note that Concavity can be proven from Subadditivity and the Chain Rule.

(c) Determine which of these properties are *also* obeyed by the Renyi entropy  $S_n = \frac{1}{1-n} \ln \text{tr}(\rho^n)$ , which limits to the von Neumann entropy as  $n \rightarrow 1$ . For simplicity you may wish to focus on the case  $n > 1$  and finite dimensional Hilbert spaces. *Note: the simplest counterexample to Concavity involves a Hilbert space with a large number of states.*