

RENORMALIZATION EXAMPLE IN $D=4$

4.5

free fields

calculate 1-loop effective action, div. pieces local:

$$I_{\text{eff}} = \int d^4x \sqrt{g} \left[\frac{R}{16\pi G} + \alpha R^2 + \beta (R_{ab})^2 + \gamma (R_{abcd})^2 \right]$$

$$\Delta \frac{1}{G} = f_G \epsilon^{-2} + I_{\text{nonlocal}}$$

$$\Delta \alpha = f_\alpha \ln(\epsilon) + g_\alpha$$

f_G, g 's nonuniversal

$$\Delta \beta = f_\beta \ln(\epsilon) + g_\beta$$

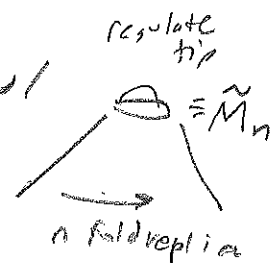
$f_{\alpha, \beta, \gamma}$ universal

$$\Delta \gamma = f_\gamma \ln(\epsilon) + g_\gamma$$

(depend on # & type of species)

FPS calculated "squashed cone" entropy associated w/ regulated replicated manifold, in $\lim n \rightarrow 1$

$$\int_{\tilde{M}_n} \sqrt{g} d^Dx R \rightarrow n \int_M \sqrt{g} d^Dx R + \underbrace{8\pi(1-n)A(\epsilon)}_{\text{from tip}} + \dots$$



$$\int_{\tilde{M}_n} \sqrt{g} d^Dx R^2 \rightarrow n \int_M \sqrt{g} d^Dx R^2 + 8\pi(1-n) \int_E d^{D-2}y \sqrt{g} R + \dots \quad \left[\mathcal{O}(1-n)^2 \right]$$

$$\int_{\tilde{M}_n} \sqrt{g} d^Dx (R_{ab})^2 \rightarrow n \int_M \boxed{\phantom{R_{ab}}} + 4\pi(1-n) \int_E d^{D-2}y \sqrt{g} (R_{ii} - \frac{1}{2} K_{ii} K_{jj})$$

$$\int_{\tilde{M}_n} \sqrt{g} d^Dx (R_{abcd})^2 \rightarrow n \int_M \boxed{\phantom{R_{abcd}}} + 8\pi(1-n) \int_E d^{D-2}y \sqrt{g} (R_{ijij} - K_{ij} K_{ij})$$

$$S_{\text{gen}} = \langle S_{\text{grav}} \rangle + S_{\text{out}}$$

$$= \lim_{\epsilon \rightarrow 0} \left[S_{\text{out}}^{(\epsilon)} + \frac{A}{4G(\epsilon)\hbar} + \frac{8\pi}{\hbar} \int_E d^2x \sqrt{g} \left[\right.$$

$$\left. \alpha(\epsilon) R + \frac{\beta(\epsilon)}{2} \left(R_i^i - \frac{1}{2} K_i^{ia} K_{ja}^j \right) + \gamma(\epsilon) \left(R_{ij}{}^{ij} - K_{ij}^a K_a^{ij} \right) \right]$$

DISCUSS CONTACT TERMS?

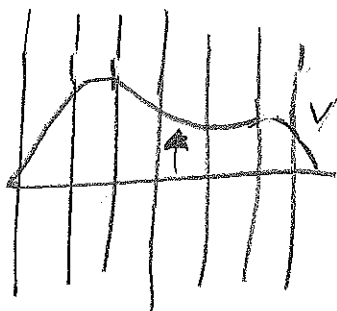
↖ where the EE formula
from earlier came from

MONOTONICITY PROPERTIES OF S_{gen}

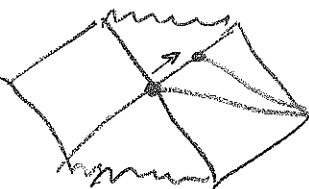
(5.1)

GSL | Increases on Causal Horizon

H



"affine"
null
Coord
v.



2nd Law requires comparing
two time slices of
horizon.

Can evolve forwards in
wiggly way.

assume bit order part. from Tab,
Semiclassical

Calculate area by focussing of light-rays from future

$$A(\infty) - A(v) = \int_v^\infty (v - v') T_{vv} dv$$

for $v' = 0$ (bit surface)
flux of Killing energy across
horizon when $v' = \text{bit surface}$

(obtained from linearized Raychaudhuri Eq. but still valid for other choices of v .)

$$\dot{\theta} = -\frac{\theta^2}{D-2} - \sigma_{ab}\sigma^{ab} - R_{ab}k^a k^b \quad ; \quad \theta = \frac{1}{A} \frac{dA}{d\lambda}$$

Use monotonicity of rel S:

$$\sigma_{ab} = \text{shear}$$

$$S(p | \sigma)_{v_2} \leq S(p | \sigma)_{v_1}$$

State
you
care about

HM
thermal
state

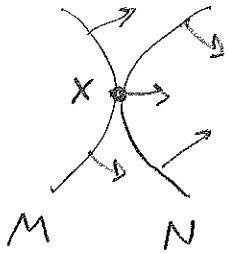
need to show state above cut v' is always thermal $e^{-2\pi K}$
where K generates dilatations about slice.

Possible to show thermality by lightfront quantization (fine print)

BEND MONOTONICITY

(S.2)

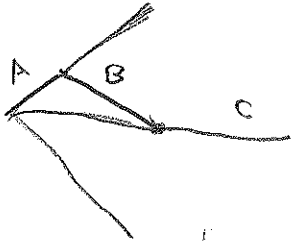
N not inside of M but touching at X



$$\left. \frac{\delta S_{\text{gen}}}{\delta V(x)} \right|_M \geq \left. \frac{\delta S_{\text{gen}}}{\delta V(x)} \right|_N$$

(at least semiclassically)

- 1) Θ piece obvious by geometry
- 2) $\frac{\delta}{\delta V} S_{\text{out}}$ piece from monotonicity of Mutual Info = SSA
Since adding extra region only makes S_{gen} increase faster.



$$S_{BC} - S_{ABC} \geq S_B - S_{AB}$$

- 3) Use $\ln G$ expansion to neglect higher cum. entropy in comparison w/ area piece.

QUANTUM FOCUSING

INFO LOSS

D:

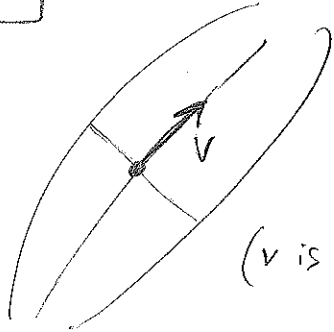
ENTROPY FOCUSES

Bousso, Fisher, Leichenauer, Wall "A quantum focussing conjecture"
 Wall "A Second Law for higher curvature gravity" + Koeller "Proof of the Quantum Null Energy Condition"

Bhattacharjee, Sarkar, Wall, "The holographic entropy increases in higher curvature gravity"

Wall "The Generalized Second Law implies a Quantum Singularity Theorem"

FOCUSING OF AREA



(v is affine)

Raychaudhuri Equation says area focusses

$$\frac{d\theta}{dv} = -\frac{\theta^2}{D-2} - \sigma_{ab}^2 - R_{\nu\nu}$$

$$\theta = \frac{1}{A} \frac{dA}{dv}$$

\Downarrow
 + Einstein $R_{\nu\nu} = 8\pi T_{\nu\nu}$
 + NEC $T_{\nu\nu} \geq 0$

$$\dot{\theta} \leq 0$$

Penrose singularity thm uses this + G.H. + space noncompact to prove that singularities must form if trapped surfaces:

compact surface whose light rays initially contract going outwards,



$$\theta \rightarrow -\infty @ v = \text{finite}$$

Also places many other important constraints on spacetimes:

- no wormholes or warp drives, AdS has boundary causality
- restrictions on baby universes, can't restart inflation
- initial singularity for noncompact FRW-like universes.

These are important results \rightarrow but do they hold behind Einstein gravity?
 \hookrightarrow like to generalize or circumvent

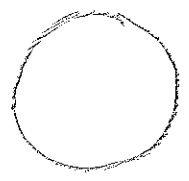
② BH entropy

Event horizons of black holes (also Rindler, de Sitter & other causal horizons) seem to obey laws of thermodynamics if you assign

$S_{BH} = \frac{A}{4G\hbar}$ (believed to count QG microstates) increases classically if NEC (Hawking)

there are corrections beyond Einstein gravity, of 2 kinds

1) Quantum/thermal corrections



$S_{out} = -tr(\rho \ln \rho)$

counts entropy of matter outside important in quantum situations e.g. Hawking radiation where sun

$\frac{\delta}{\delta V} S_{gen} \geq 0$

$S_{gen} = S_{BH} + S_{out}$ is increasing (GSL)

S_{out} includes a piece coming from entanglement entropy near horizon, which is divergent & must be renormalized

proven semiclassically Wall II for free/superren. fields minimally coupled to Einstein gravity & mod. or rel S

2) higher curvature (counterterms / stringy) corrections

$I = \int d^D x \sqrt{-g} L, \quad L = \frac{R}{16\pi G} + \alpha(R^2 \dots)$

corrections to BH entropy

$S_{BH} = \underbrace{-\frac{2\pi}{\hbar} \int d^{D-2} x \sqrt{g}}_{\text{Wald, valid stationary, } f(R)} + 4 \frac{\partial L}{\partial R_{\mu\nu\mu\nu}} + 16 \frac{\partial^2 L}{\partial R_{\mu\nu\mu\nu} \partial R_{\lambda\rho\lambda\rho}} K_{\mu\nu}^{(\lambda\rho)} K_{\lambda\rho}^{(\mu\nu)} + \mathcal{O}(R^3) \dots$

divergences from S_{out} renormalize $\frac{1}{G}$, etc in a consistent way Solodukhin, FPS, Dong, Camps, Miao, f(Riemann)

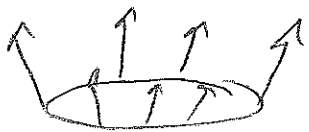
③ Spacetime thermodynamics

Evaluate S_{gen} for non-horizon slices
 (must be codim 2 split of Cauchy slice)

Research program: extend classical GR results to
 semiclassical gravity by replacing A w/ S_{gen} everywhere
 (Wall 11b, FLMB, Engelhardt + Wall 14, Bousso-Engelhardt...)

GSL + G.H. \Rightarrow quantum singularity theorems
 (w/o NEC) \Rightarrow (eg. at least semiclassically,
 spatially noncompact bounces impossible)

refer back to list



"quantum trapped surface"

$$\ominus \equiv \frac{\delta S_{gen}}{\delta V} = \frac{1}{A} \frac{\partial S_{gen}}{\partial V} \leq 0 \text{ everywhere}$$

proof rather global...

is there analogue of focussing thm?

$$\text{QFC: } \frac{\delta \Theta}{\delta V} \leq 0 \quad (\text{related to Bousso bound})$$

will give partial proofs, in semiclassical regime (t_h, α expansions);
 corrections important only if null surface is quasi-stationary,

α must be small
 "string" or pluck
 scale for thm.
 to be causal
 (Camanho et al. 14)

or else classical Einstein focussing dominates.
 (for "adiabatic processes", $\delta S_{gen} \approx 0$ and all terms matter)

MAYBE NOT SAY:

$A/\mu\text{gth}$ term may still refer to Θ & microstates.

(Bianchi-Myers 12) or something more subtle like diff. ent.

(4) Higher curvature corrections

comment?

Start w/ covariant action with finite α corrections

$$I = I_{\text{grav}}[g^{ab}, R_{abcd}; \nabla R \dots \phi, \nabla \phi] + I_{\text{mat}}[g^{ab}, \psi]$$

↓
arbitrary h.c. scalar-tensor
action for grav. fields

↓
assume minimal,
NEC

EOM $H_{ab} = T_{ab}$, $H_{kk} = T_{kk} \geq 0$

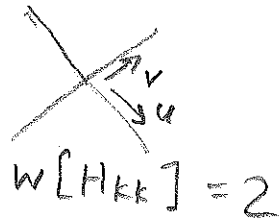
assume g_{ab}, ϕ are first order perturbations to a regular bifurcate Killing horizon [a bit stronger than $\Theta_{cl} = 0$]

▷ if u, v are the two null directions,

any quantity w / $w = \#v\text{'s} - \#u\text{'s}$

scales like v^{-w} on Killing spacetime,

if $w > 0$, must vanish to be regular @ $v=0$
at 0th order



Using this fact, can show that:

$$\delta H_{vv} = \sum_n X^{(-n)} \delta Y^{(2+n)} \Rightarrow \int_H^{D-2} \sqrt{g} T_{vv} = \frac{\hbar}{2\pi} \partial_v \partial_v \xi$$

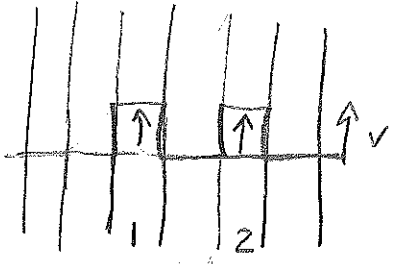
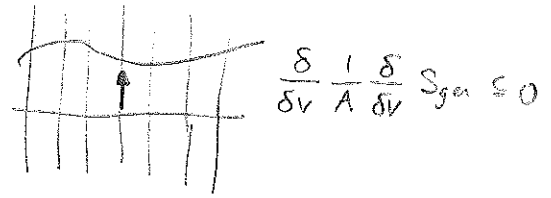
- gauge fix metric so $g_{vv}, g_{vi}, g_{vu}, g_{vu} = 0$ so that $w=2$ implies at least 2 ∂_v 's.

- repeatedly diff. by parts until 2 ∂_v 's are outside expression,

using $\partial_v[w \geq 0] \partial_v[w \geq 0] = 0$ @ 1st order \rightarrow only δ term has $w \geq 0$

can show $\xi = \delta S$, for f(Riemann), $S = S_{\text{Dong}} + \text{total deriv.} + \mathcal{O}(k^4)$
implies S_{Dong} obeys 2nd Law.

(5) Quantum focussing



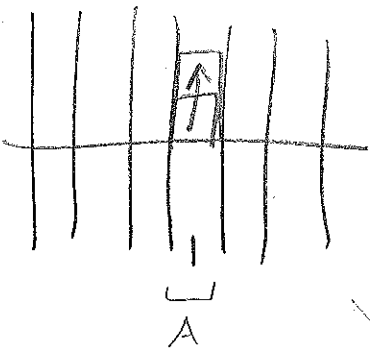
Can also consider effects of quantum fields falling across a stationary null surface - discretize into pencils

divide into mixed case

$$\frac{\partial}{\partial v_1} \frac{\partial}{\partial v_2} S_{out} \leq 0 \text{ by strong subadditivity}$$

area does not contribute

or into "contact term"



A

QNEC

$$\frac{\delta^2 S_{out}}{\delta v(y) \delta v(y')} = f(y) \delta(y-y') + \text{off-diagonal}$$

$$\frac{\partial}{\partial v_1} \frac{\partial}{\partial v_1} S_{gen} \leq 0 \Rightarrow \langle T_{vv} \rangle \geq \frac{\hbar}{2\pi A} S''_{out}$$

Stronger than SSA

restrict to 2d field theory near the vacuum (UV limit)

$$P = P_0 + \delta P$$

2nd order in δP , can be entangled w/ auxiliary system

Calculated entropy using replica trick and found it was true



- Continue to noninteger n
- $S = (1-n \partial_n) \ln Z$
- Evaluate @ n=1

New law of quantum information theory?

S_{gen} closely related to "relative entropy" (Casini)

But doesn't follow from standard properties like positivity or monotonicity since it involves a 2nd time derivative